



Spin-Orbit Coupling in SIESTA

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First steps with SIESTA from zero to hero

siesta

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Spin-Orbit Coupling



Spin-orbit coupling (SOC)

• relativistic effect interaction of the electron spin \widehat{S} and its own orbital momentum \widehat{L}

 $\widehat{H}_{SO} \sim \widehat{S} \cdot \widehat{L}$

Rest frame of the nucleus



Rest frame of the electron



Isolated atom:

split energy levels in atomic fine structure

• General:

couples real space and spin space

 \rightarrow also spins & crystal lattice

- Nucleus motion creates magnetic field $\widehat{B} \sim \widehat{L}$
- Electron spin couples to \widehat{B}
 - \rightarrow Zeemann-like correction: $\hat{H}_{SO} \sim \hat{\mu} \cdot \hat{B} \sim \hat{S} \cdot \hat{L}$
- without SOC:

(*l*, *m*) and spin (*s*) are good quantum numbers

• with SOC:

use (l_J, m_J) quantum numbers related to $\hat{J} = \hat{L} + \hat{S}$

Spin-Orbit Coupling



SOC: small correction leading to rich physics

- Giant Magneto Resistance (Nobel prize 2007)
 → read heads for hard drives
- Spin transfer torque

 → electrically changing the magnetic configuration of a device
- Topological insulators
- ... and many more ...

So far: SIESTA only without spin

- How do we include spins?
- How do we include spin-orbit coupling?



Manchon, Nature Mater 14 (2015).

Spins in DFT & SIESTA



How can we control the spins in SIESTA?

Flag Spin

- non-polarized: no spins (default)
- **polarized**: colinear spins
- **non-colinear**: non-colinear spins no spin-orbit coupling
- **spin-orbit**: non-colinear spins with spin-orbit coupling



non-colinear



What does it mean?

Two Hamiltonians $\widehat{H} \to \widehat{H}_{\uparrow}$, \widehat{H}_{\downarrow} Two sets of wavefunctions $|\psi\rangle \to |\psi_{\uparrow}\rangle$, $|\psi_{\downarrow}\rangle$ Two density matrices $\widehat{\rho} \to \widehat{\rho}_{\uparrow}$, $\widehat{\rho}_{\downarrow}$

Spin-channels fully decoupled? No!

- Only in the Schrödinger equation.
- Exchange-correlation functional couples both spin-channels:

```
E_{XC}[\rho] \rightarrow E_{XC}[\rho_{\uparrow},\rho_{\downarrow}] \neq E_{XC}[\rho_{\uparrow}] + E_{XC}[\rho_{\downarrow}]
```

Colinear spin calculation:

no SOC

 \rightarrow spin and real space decoupled

 \rightarrow polarization axis: arbitrary in SIESTA always **z-axis**



- Information printed each SCF step:
 - includes (x,y,z) component of the spin moment of the whole cell and total spin.
 - in collinear case x and y components always zero.

	isc	f Eharris(eV)	E_KS(eV)	FreeEng(eV)	dDmax	Ef(eV)	dHmax(eV)	
scf:		1 -2153.25	7209	-2198.764140	-2198.764152	3.041441	-3.686948	23.600408	
sp:	in m	<pre>oment: {S}</pre>	, S	= { 0.0	0.0	2.53181 }	2.53181		



- Information printed each SCF step:
 - includes spin moment and (x,y,z) component.
 - in collinear case x and y components always zero.
- .bands file: twice as many eigenvalues for each k point first channel 1 then channel 2

-3.5822	8951613775	5							
0.0000	0000000000)E+000 4	i.14093362	787451					
-13.841	9785400988	3 6	8.6417154	298812					
	24	2	351			# Number	of spin c	hannels 2	
0.00000	-13.8420	-10.1800	-9.8289	-7.3976	-7.3976	-6.7511	-6.0972	-5.8342	-4.5706
	-4.4758	5.2572	19.1906	19.5233	27.6085	27.6085	31.6350	34.5379	37.6186
	50.6420	56.7901	61.3181	61.3181	-13.5977	-9.5698	-9.4133	-6.2825	-6.2825
	-4.5791	-3.2974	-2.1377	-2.1377	-2.1064	6.2058	20.0516	20.3546	28.3583
	32.2540	35.4856	38.8661	43.4723	52.7600	57.5402	62.4527	62.4527	
0.008768	-13.8407	-10.1792	-9.8285	-7.3971	-7.3971	-6.7516	-6.0970	-5.8347	-4.5708



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- gnubands:
 - use option **-s** to select spin up (1) or spin down (2)
 - or nothing: adds third column to output with 1 and 2 to identify spin up and spin down





How do SIESTA outputs change?

- Information printed each SCF step:
 - includes spin moment and (x,y,z) components.
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• DOS file:

- one additional column per energy point: energy, DOS(up), DOS(down)



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 - one additional column per energy point: energy, DOS(up), DOS(down)
- ... quantities separately for up and down channel

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non-colinear

Spinor wavefunction $|\psi\rangle \rightarrow \begin{pmatrix} |\psi_{\uparrow}\rangle \\ |\psi_{\downarrow}\rangle \end{pmatrix}$

Spin dependent operators

$$\widehat{H} \to \begin{pmatrix} \widehat{H}_{\uparrow\uparrow} & \widehat{H}_{\uparrow\downarrow} \\ \widehat{H}_{\downarrow\uparrow} & \widehat{H}_{\downarrow\downarrow} \end{pmatrix} \qquad \widehat{\rho} \to \begin{pmatrix} \widehat{\rho}_{\uparrow\uparrow} & \widehat{\rho}_{\uparrow\downarrow} \\ \widehat{\rho}_{\downarrow\uparrow} & \widehat{\rho}_{\downarrow\downarrow} \end{pmatrix}$$

What do the values of $\widehat{
ho}_{\sigma\sigma'}$ indicate?

- $\widehat{
 ho}_{\uparrow\uparrow}+\widehat{
 ho}_{\downarrow\downarrow}$: related to charge
- $\widehat{\rho}_{\uparrow\uparrow} \widehat{\rho}_{\downarrow\downarrow}$: related to spin polarization *z*-axis
- Re{ $\hat{\rho}_{\uparrow\downarrow}$ }: related to spin polarization along *x* axis
- Im{ $\hat{\rho}_{\uparrow\downarrow}$ }: related to spin polarization along *y*-axis

 $Q = Tr\{ (\hat{\rho}_{\uparrow\uparrow} + \hat{\rho}_{\downarrow\downarrow}) S \}$ $M_{z} = \mu_{B} \frac{1}{2} Tr\{ (\hat{\rho}_{\uparrow\uparrow} - \hat{\rho}_{\downarrow\downarrow}) S \}$ $M_{x} = \mu_{B} Tr\{ Re\{ \hat{\rho}_{\uparrow\downarrow} \} S \}$ $M_{y} = \mu_{B} Tr\{ Im\{ \hat{\rho}_{\uparrow\downarrow} \} S \}$





How can $\widehat{H}_{\uparrow\downarrow}$ be non-zero?

1. If not all spins are align along z, then

$$\widehat{
ho}_{\uparrow\downarrow}
eq 0
ightarrow$$

To calculate E_{XC}

a) the density is locally (i.e. for each point on the grid) rotated along z

b) $E_{XC}[\mathbf{\rho}_{\uparrow}, \mathbf{\rho}_{\downarrow}]$ is calculated

c) the corresponding term in the Hamiltonian is rotated back according to original direction

 $\rightarrow \widehat{H}_{\uparrow\downarrow} \neq 0$

2. If we include spin-orbit coupling:

 $\widehat{H}_{SO} \sim \widehat{L} \cdot \widehat{S} \\ \rightarrow \widehat{H}_{\uparrow\downarrow} \neq \mathbf{0}$



- Information printed each SCF step:
 - includes (x,y,z) component of the spin moment of the whole cell and total spin.
 - now all components can be non-zero

	iscf	Eharris(eV)	E_KS(eV)	FreeEng(eV)	dDmax	Ef(eV)	dHmax(eV)
scf:	1	-2153.257209	-2198.764140	-2198.764152	3.041441	-3.686948	23.600408
sp	in mo	<pre>ment: {S} , S </pre>	= { 2.53181	2.53181	2.53181 }	4.38522	



- Information printed each SCF step:
 - includes spin moment and (x,y,z) component.
 - now all components can be non-zero
- .bands file: twice as many eigenvalues for each k point, no separation by spin index

-3.5822	8951613775 00000000000	5)E+000 4	4.14093362	787451					
-13.041	7/03400700		DO.041/104	270012					
	48	1	351						
0.00000	-13.8461	-13.6001	-10.2404	-9.9188	-9.6644	-9.5335	-7.4540	-7.4037	-6.8073
	-6.3620	-6.0184	-5.7102	-4.8477	-4.7915	-4.4688	-4.3186	-4.2566	-3.2346
	-2.0978	-2.0153	5.2669	6.2161	19.1865	19.5289	20.0456	20.3643	27.5877
	28.3833	28.3866	31.6409	32.2675	34.5233	35.4643	37.6469	38.8891	42.3467
	50.6597	52.7767	56.7916	57.5425	61.3260	61.3262	62.4743	62.4752	
0.008768	-13.8448	-13.5988	-10.2396	-9.9184	-9.6636	-9.5332	-7.4535	-7.4032	-6.8076



- Information printed each SCF step:
 - includes spin moment and (x,y,z) component.
 - in collinear case x and y components always zero.
- .bands file: twice as many eigenvalues for each k point, no separation by spin index
- **gnubands**: like no-spin-case, but twice as many bands

0.000000	-10.259710
0.008768	-10.258410
0.017536	-10.254510
• • •	
0.000000	-10.259710
0.008768	-10.258410
0.017536	-10.254510
• • •	



How do SIESTA outputs change?

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- .bands file: twice as many eigen values for each point first channel 1 then channel 2
- gnubands:
 - use option **-s** to select spin up (1) or spin down (2)
 - or nothing: adds third column to output with 1 and 2 to identify spin up and spin down

• DOS file:

- one additional column per energy point: energy, DOS(up), DOS(down)

-9.99999	1.00971	0.63864
-9.98497	1.04733	0.66864
-9.96996	1.04293	0.64380
	1.04270	



- Information printed each SCF step:
 - includes spin moment and (x,y,z) component.
 - in collinear case x and y components always zero.
- .bands file: twice as many eigenvalues for each k point, no separation by spin index
- **gnubands**: like no-spin-case, but twice as many bands
- DOS file:
 - two additional columns: *energy*, *DOS*(↑↑), *DOS*(↓↓), *Re{DOS*(↑↓)}, *Im{DOS*(↑↓)}

Spin Orbit: Example 1



Siesta Documentation

0.1

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Molecular Dynamics

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Polarization calculations with the Berry-phase approach

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Advanced analysis of the electronic structure

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Spin-Orbit coupling

Author Ramón Cuadrado (University of Southampton) & Nils Wittemeier (ICN2)

The main goal of this tutorial is to cover the basic notions to perform calculations using the spin and spin-orbit coupling (SOC) in Siesta.

Performing calculations with different spin configuration.

Hint

Please move to directory Example-1.

Note

The original fdf for this example can be found in *example_1.fdf.original*. We have made some changes (can you see them?) to speed the calculation while maintaining reasonable quality results.

In this example you will calculate the band structure, density-of-states (DOS), and Mulliken charges of bulk FePt using different spin options.

Getting started

To get started inspect the input file *example_1.fdf* file, and make sure that you are familiar with all options in the input file. Refer to the user manual to clarify any doubts. Alongside the input file you will find two pseudopotential files one for Fe and one for Pt (*Fe_fept_SOC.psf*, *Pt_fept_SOC.psf*)



Perform and analyze calculations of FePt with all different option for *Spin*

Spin-Orbit coupling in SIESTA



How can we include SO in SIESTA? - Pseudo potentials.

Pseudo potential generation

Scalar relativistic case:

- 1. Solve **Schrödinger** equation including scalar relativistic corrections
- 2. Smoothen wavefunction
- 3. Construct potential

$$\hat{V}_{PP} = V_{\text{local}}(r) + \sum_{l} \delta V_{l}(r) \hat{P}_{l}$$

Fully relativistic pseudo potential:

- 1. Solve **Dirac** equation
- 2. Smoothen wavefunction
- 3. Construct potential

•
$$\mathbf{j} = \mathbf{l} \pm \frac{1}{2}$$
 instead \mathbf{l} $(\hat{\mathbf{j}} = \hat{\mathbf{L}} + \hat{\mathbf{S}})$
 $\hat{V}_{PP} = V_{\text{local}}(r) + \sum_{j} \delta V_{j}(r)\hat{P}_{j}$

• Twice the number projectors and PP components

R Cuadrado and J I Cerdá 2012 J. Phys.: Condens. Matter 24 086005

Spin-Orbit coupling in SIESTA



How can we include SO in SIESTA? - Pseudo potentials.



R Cuadrado and J I Cerdá 2012 J. Phys.: Condens. Matter 24 086005

Spin-Orbit in SIESTA



How to run SIESTA with spin-orbit coupling

- use spin spin-orbit
- use fully relativistic pseudopotentials
- carefully converge
 - k-points sampling & electronic temperature
 - mesh cutoff
- Use low tolerance for the
 - density matrix $< 10^{-5} \text{ eV}$
 - And/or Hamiltonian $< 10^{-4} \text{ eV}$

Specific outputs:

- Energy decomposition:
 - Eso: spin-orbit energy
- Everything else like non-collinear case



Example: Magnetic Anisotropy (CoPt nano particle)



 \mathcal{Y}

•

•

Example: Magnetic Anisotropy (CoPt nano particle)

(001) (010)35 30 [meV] 25 (100 L° Etot. 125 150 175 100 θ (degree)



As the angle between the magnetic

When the spin is (anti)parallel to

the z-axis the energy is the lowest,

Magnetic anisotropy energy (MAE):

 $E_{tot}(\theta=0) - E_{tot}(\theta=90) \sim 36 \text{ meV}$

i.e. this is the **most stable**

total energy rises.

configuration.

moment and the z-axis increases, the







 $x_{\mathcal{E}}$

Example: Magnetic Anisotropy (CoPt nano particle)

Disclaimer

During the SCF cycle, the magnetic moments may change. **Spin directions are not fixed.** As such the final and initial spin moments may differ. The final moments should converge to a local minimum on the energy surface. In practice, the spin moments often do not fully converge and it is possible to explore different spin configurations. However, **always check the final spin moments!**

Ideally, one should use **non-colinear spin-constrained DFT*** calculations. This feature was already implemented in SIESTA but is not publicly available, yet.

*R Cuadrado et al 2018 J. Phys. Mater. 1 015010

Spin Orbit: Examples 2 and 3

EXCELENCIA SEVERO OCHOA

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Calculation of optical properties

Computing magnetic interaction parameters with TB2J

Simulation of STM images

Calculation of the magnetic anisotropy of Pt₂.

In this example, you will calculate the magnetic anisotropy of a Pt_2 dimer. This means you will have to run multiple SIESTA calculations each with magnetic moments aligned along different directions and extract the total energy.

To get started inspect the input file *example_2z.fdf* file. Use this input file and the pseudopotential of Pt atom, $Pt_pt2_SOC.psf$ (note that the pseudo for Pt is different from the one used in the first example), and calculate the total self-consistent energy Etot for three highly symmetric orientations of the initial magnetization, namely, X, Y and Z, keeping fixed the physical dimer (the (x,y,z) coordinates of each atom).

Each calculation has to be initialized modifying the DM.InitSpin, updating the orientation of each atomic magnetization. As a result one has to have three different energy values for each (θ, ϕ) pair.

Remember to rename each one of the .*fdf* files, change the SystemLabel flag inside for each magnetic orientation and do not re-use the density matrix (*.DM* file) obtained in other orientations. Reading the density matrix will cause SIESTA to ignore the <u>DM.InitSpin</u>. To calculate the total energy for each specific magnetic configuration you need to start from scratch.

Calculation of the magnetic anisotropy of FePt bulk

In this example, you will calculate the magnetic anisotropy bulk FePt.

You will find the *example_3z.fdf* file and the PPs of Fe and Pt: $Fe_fept_SOC.psf$ and $Pt_fept_SOC.psf$. Each calculation has to be initialized modifying the <code>DH.InitSpin</code> block, as above, updating the orientation of each atomic magnetization. As a result one has to have different energy values for each (θ, ϕ) pair.

• A) Obtain the Etot for three highly symmetric orientations of the initial magnetization, namely, X, Y and Z, keeping fixed the atoms in their unit cell. Again do not use the same density matrix for same calculations. Allow SIESTA to calculate it from scratch.

• B) Create a magnetization curve changing the angles at intervals of 20 degrees from the Z-axis initial orientation, to finalize along the X axis. Plot the resulting energy as a function of the varied angle. Check the orientation of the final magnetic moment in each output file. Does it coincide with the initial one?

- In addition to the first exercise, We prepared two more exercise, in which we walk through the calculation of the magnetic anisotropy of
 - a Pt₂ dimer
 - bulk FePt

